Influence of Glass Composition on the Luminescence Mechanisms of CdSe Quantum-Dot-Doped Glasses

Wenke Li, Xiujian Zhao, Chao Liu,* and François-Xavier Coudert*

Abstract: In this work, we characterized the electronic structure of CdSe quantum dots embedded in a series of x Na2O, (1−x) SiO2 glass matrices (x = 0, 0.25, 0.33, and 0.5). We analyzed the impact of the glass matrix composition on both the atomic structure of the quantum dot (QD) and the QD/glass interface, as well as the luminescence mechanisms, using density functional theory calculations. The increase of Na2O content in the glass matrices was found to promote the formation of Cd–O and Se–Na interfacial bonds and disrupt the Cd–Se bonds network. In particular, we show that the glass composition directly affects the nature of the highest occupied molecular orbitals (HOMO). According to the atomic structure, band gap distribution, and density of states calculation, we find that there is significant reconstruction of the QD and that the picture sometimes proposed of a “pristine quantum dot” surrounded by glass is not realistic. The introduction of CdSe QD significantly decreased the HOMO–LUMO gap of the glass compared to pristine glasses, and the interfacial bonds greatly contributed to the frontier orbitals without forming midgap states. We propose a new energy diagram, quite different from the traditional model, to explain the luminescence of CdSe quantum-dot-doped glasses, originating from the intrinsic emission of this hybrid system {QD + glass}. These results improve our understanding of the luminescence of CdSe quantum-dot-doped glasses, explaining the reason for the poor quantum efficiency and broad emission linewidth compared with their colloidal counterparts.

Introduction

Quantum dots (QDs) exhibit remarkable electronic and optical properties induced by quantum confinement. Colloidal quantum dots can be produced at a low cost with a high quantum efficiency and narrow size distribution, which makes them promising for photovoltaics,1 lighting,2 and labeling3 applications. Incorporation of QDs into glasses can prevent the agglomeration and enhance the chemical and thermal abilities, making them appealing for nonlinear optical devices and LEDs.4

Due to the high surface-to-volume ratio of quantum dots, the main obstacle for the success of their applications is to achieve a precise control of the emission properties, which mainly depends on the surface chemistry of the nanocrystals, and the interface with the surrounding medium. Surface defects are ubiquitous to QDs, and an in-depth understanding of their atomic origin is necessary to tune the emission and design defect-free QDs. In spite of a wealth of experimental data reported for CdSe QDs, the question of exactly how surface defects influence the photophysical properties of QDs is still open. It is experimentally challenging to resolve spectral features originating from QD–ligand interactions. Density functional theory (DFT) is able to address these issues, gaining comparable gaps between the highest occupied (HOMO) and the lowest unoccupied molecular orbitals (LUMO), IR, and Raman spectra to experimental findings.6

In the literature, Cd33Se33 QDs are the most commonly used representative models of quantum dots for computational studies, as they are believed to have the highest degree of stability, exhibiting basic optical features similar to larger QDs. Kilina et al.7 found that the electron trap states were very sensitive to ligand position and the solvent polarity. Sargent et al.8 demonstrated that surface vacancies can improve the fluorescence yield compared to vacancy-free surfaces. Beratan et al.9 performed comparative studies of the structural and electronic properties of pristine and NH3−, SCH−, and OPH− capped Cd33Se33 QDs, showing that an increase in capping ligand concentration causes noticeable changes in the capped QD electronic structure, leading to different surface states and HOMO–LUMO gaps. Pudzder et al.10 reported significant reconstruction of the nanoparticle surface, while the wurtzite core was maintained, leading to the opening of an optical gap without the aid of passivating ligands, thus “self-healing” the surface electronic structure.

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All of those theoretical studies helped understand the structural origin of defect emission and effective surface passivation of colloidal QDs. However, relatively little research effort has been devoted to the computational study of CdSe quantum dots embedded in glasses compared to their colloidal counterparts. The quantum efficiency of the quantum dots embedded in glass matrices is extremely low, which is usually attributed to their surface passivation and broad size distribution. Recent work showed that the size distribution can be controlled by tuning the heat treatment and the concentration of dopants. However, surface control techniques used in colloidal QDs, such as post-synthetic passivation by surfactants or growth of an iso-structural shell, are difficult to perform when the QDs are surrounded by dense and amorphous glass matrices.

Hence, defect emission, rather than intrinsic emission, is most frequently observed in the CdSe quantum-dot-doped glasses. It is of predominant importance to fully comprehend the structural origin of defect emission and the impact of QD/glass interface on the electronic structure. The atomic structure of the CdSe quantum-dot-doped glasses is exceedingly difficult to characterize due to the typical low density of QDs in the glass and the instability of glass matrices under electron beam. In a previous study, we used network modifiers as well as nonbridging oxygen atoms to be capped with Cd33Se33 clusters, to explore the impact of these additional surface atoms on the morphology and electronic structure on the CdSe QD. However, this simplistic model was far from being a realistic environment of the quantum dot in glass matrices. In later work, we created a more detailed atomic structure using a combination of classical and ab initio molecular dynamics, demonstrating the complex interfacial chemical environment between the CdSe quantum dot and a surrounding glass matrix.

In the present work, we used DFT-based methods to determine the structure of CdSe quantum-dot-doped glasses with varying composition and to determine their electronic structure to probe the impact of the glass matrices on the optical properties of the CdSe quantum dot. Besides, the sodium ion was found to sharply decrease the HOMO–LUMO gap in previous DFT calculations. In their role of glass modifiers, sodium ions can alter the structure of silicate glasses, and it was experimentally observed that there is a disappearance of the visible absorption of Se−Se color centers upon increasing concentration of Na2O in a silicate glass doped with ZnSe. The Se−Se color centers are considered to be the nucleation sites of the CdSe QD. Therefore, in this study, we changed the amount of Na2O in the glass (Na2O)x(SiO2)1−x (x = 0, 0.25, 0.33, and 0.5 in molar fraction) to explore the compositional dependence of the atomic and electronic structures of CdSe quantum dot in glass matrices. The results demonstrate that an increase in the amount of Na2O contributes to the formation of Cd−O and Se−Na bonds, with breaking of Si−O, O−Na, and Cd−Se bonds. The density of states (DOS) and projected density of states (PDOS) were also calculated. Although with the same composition, the dominant luminescence mechanisms in different configurations are also different. The top of the valence band and the bottom of the conduction band are decided by the hybrid QD in most compositions. However, in most configurations of CdSe, quantum-dot-doped glass with composition 0.33 Na2O−0.67 SiO2 exhibits totally different luminescence mechanisms that the top of the valence band is determined by the hybrid glasses.

Most importantly, based on the analysis of atomic structure, HOMO–LUMO gap distribution, and DOS, the intrinsic emission from pristine QD is found to be negligible due to the little possibility of QD without any interaction with glass matrix. The photoluminescence of CdSe quantum dots doped glass is originated from the intrinsic emission of these complicated hybrid system rather than the intrinsic or defect emission of pristine QD, which is a common model proposed by experimentalists. These results provide a better understanding of the electronic structure and luminescence mechanisms of these complex systems, giving guidance for future compositional design of highly luminescent glass containing quantum dots.

### COMPUTATIONAL METHODS

We generated configurations of pristine sodium silicate glasses and, from those pristine glasses, created CdSe quantum-dot-doped silicate glasses, using a combination of classical force-field-based molecular dynamics (MD) and ab initio molecular dynamics (AIMD). The protocol followed here is adapted from that developed and validated in our previous work. The pristine QD was cut from the bulk structure and then relaxed to the lowest energy configuration. The detailed methodology involved in the generation of glass matrix and CdSe quantum-dot-doped glasses is presented in the Supporting Information.

For each system, 500 configurations were selected at equal intervals from a 10 ps production run of ab initio molecular dynamics simulations. Among those, 50 configurations were picked to perform the calculation and analysis of the electronic structure of CdSe quantum-dot-doped glasses. We used the nonlocal PBE0-TC-LRC exchange-correlation functional with a cutoff radius of 2.0 Å. From previous work, the electronic structure calculations based on the generalized gradient approximation (GGA) or local density approximation (LDA) exchange-correlation functionals severely underestimated the HOMO–LUMO gap and long-range Coulombic interactions. The inclusion of the Hartree–Fock exchanges was found in the literature, to provide a more accurate description of the band gap. Moreover, the computational cost of nonlocal functional calculations can be reduced using the auxiliary density matrix method (ADMM) (Table 1).

### RESULTS AND DISCUSSION

The final structures of the CdSe quantum-dot-doped glasses with different chemical composition (Na2O:SiO2 ratio) are depicted in Figure 1. We provide a quantitative analysis of the chemical environment of each structure in Table 2. As glass modifiers, sodium ions can break the Si−O bond between Si

### Table 1. Composition of Different x Na2O−(1−x) SiO2 (in Molar Fraction) Glasses Investigated in This Study and Their Respective Densities*

<table>
<thead>
<tr>
<th>(Na2O)x(SiO2)1−x</th>
<th>density (g/cm³)</th>
<th>O</th>
<th>Si</th>
<th>Na</th>
<th>Cd</th>
<th>Se</th>
</tr>
</thead>
<tbody>
<tr>
<td>x = 0</td>
<td>2.20</td>
<td>224</td>
<td>112</td>
<td>33</td>
<td>33</td>
<td></td>
</tr>
<tr>
<td>x = 0.25</td>
<td>2.43</td>
<td>210</td>
<td>90</td>
<td>60</td>
<td>33</td>
<td></td>
</tr>
<tr>
<td>x = 0.33</td>
<td>2.49</td>
<td>195</td>
<td>78</td>
<td>78</td>
<td>33</td>
<td></td>
</tr>
<tr>
<td>x = 0.5</td>
<td>2.56</td>
<td>186</td>
<td>62</td>
<td>124</td>
<td>33</td>
<td></td>
</tr>
</tbody>
</table>

*For each system, we list the number of O, Si, Na, Cd, and Se atoms in the simulation cell of the CdSe quantum-dot-doped glass.
and O ions, contributing to the appearance of nonbridging oxygen atoms in the inorganic network. The percentage of O atoms bonded with Si atoms decreased from 96.8 to 32.0% upon an increase in the amount of Na2O in the glass. Meanwhile, the nonbridging oxygen preferred to be bonded with Cd atoms rather than Na atoms, as demonstrated by an increase in the percentage of O atoms bonded with Cd atoms from 2.9 to 5.9%. Se atoms also tended to be bonded preferentially with sodium ions as the amount of sodium oxygen becomes larger in the glass matrices, with 37.1% of Se atoms bonded with Na atoms when \( x = 0.5 \). In some configurations, up to 33 Se atoms were found to be bonded with sodium ions, an increased breakage of the bond between Cd atoms and Se atoms when adding sodium oxygen into the glass matrices was observed.

These results showed that, near the quantum dot/glass interface, the CdSe QD is quite sensitive to the composition of the glass matrix. This has been experimentally suspected based on the Raman spectroscopic data, showing that the Raman shifts of CdSe QDs with identical size encompassed in different host matrices were different. In previous work, a lattice contraction effect, depending on the thermodynamic conditions decided by the interactions between QDs and glass matrix, was proposed to explain this observation—but without further direct evidence. In the colloidal method, the CdSe QD core was observed to maintain its bulk structure when capped by organic ligands. However, the interfacial structure and the structure of the CdSe QD reconstructed enormously in our results.

First of all, we want to outline some definitions here to better illustrate the complexity of the hybrid systems considered in this work. We divide each configuration of CdSe quantum-dot-doped glasses into three regions of space: (a) the hybrid QD; (b) the hybrid glasses, in the interfacial region; and (c) the pristine glass, at a distance such that its structure is not affected by the QD. When the pristine CdSe QD is introduced into the glass matrix, an important structural reconstruction is observed, and we note the resulting structure as a “hybrid QD” (including Cd and Se atoms) to highlight that its structure is very different from the pristine QD. At the same time, the structural reconstruction also happens in the glass matrix due to interactions with QD, forming a “hybrid glass” (including O, Na, and Si atoms). The hybrid QD and hybrid glass are connected by forming interfacial bonds such as Na–Se and Cd–O linkages. However, the density of the QDs in the glass matrix is small;

Table 2. Average Percentage of O, Na, Si, Se, and Cd Bonded with Different Atoms Averaged over the AIMD Simulations of CdSe QD in \( x \) Na2O–(1–\( x \)) SiO2 Glass Matrix

<table>
<thead>
<tr>
<th>bonds</th>
<th>( x = 0 )</th>
<th>( x = 0.25 )</th>
<th>( x = 0.33 )</th>
<th>( x = 0.5 )</th>
</tr>
</thead>
<tbody>
<tr>
<td>O–Si</td>
<td>96.79</td>
<td>59.44</td>
<td>50.29</td>
<td>31.97</td>
</tr>
<tr>
<td>O–Na</td>
<td>36.82</td>
<td>44.54</td>
<td>62.16</td>
<td></td>
</tr>
<tr>
<td>O–Cd</td>
<td>2.91</td>
<td>3.72</td>
<td>5.17</td>
<td>5.87</td>
</tr>
<tr>
<td>Si–O</td>
<td>97.99</td>
<td>96.25</td>
<td>94.54</td>
<td>83.07</td>
</tr>
<tr>
<td>Na–O</td>
<td>87.45</td>
<td>86.21</td>
<td>83.57</td>
<td></td>
</tr>
<tr>
<td>Na–Se</td>
<td>7.40</td>
<td>8.05</td>
<td>7.73</td>
<td></td>
</tr>
<tr>
<td>Se–Na</td>
<td>17.06</td>
<td>23.72</td>
<td>37.09</td>
<td></td>
</tr>
<tr>
<td>Se–Cd</td>
<td>89.61</td>
<td>81.90</td>
<td>76.26</td>
<td>62.88</td>
</tr>
<tr>
<td>Cd–O</td>
<td>12.88</td>
<td>19.87</td>
<td>27.70</td>
<td>35.95</td>
</tr>
<tr>
<td>Cd–Se</td>
<td>87.10</td>
<td>79.89</td>
<td>71.72</td>
<td>59.73</td>
</tr>
</tbody>
</table>
thus, the incorporation of CdSe QD can only influence the
glass structure at a short range, without further impact on the
glass structure in the long range. Therefore, the structure of the
pristine glass, which is far from the CdSe QD, remains
unchanged. These micro regions are separated by dense and
amorphous pristine glass.

Within our chosen methodology, the HOMO−LUMO gap
for the pristine QD is found to be 2.08 eV, which is smaller
than the experimental value (∼3 eV). The goal of our
present work is not to reproduce the exact value of the
experimentally observed absorption peak, but rather to predict
the impact of the glass matrices on a relative energy scale,
which can be corrected “a posteriori” by a constant energy shift
to a first approximation. To perform significant statistical
sampling of different configurations obtained from AIMD, we
are also constrained by computational cost. In addition, 10
configurations were selected from the AIMD production run
of pristine glasses, and the average HOMO−LUMO gaps are
calculated to be 7.19, 4.44, 4.33, and 4.41 eV when
x = 0, 0.25, 0.33, and 0.5, respectively. The reported experimental value for
amorphous SiO₂ lies in the range of 8.7−9.4 eV. The
calculated HOMO−LUMO gap of SiO₂ glass in our work was
comparable to other researchers’ work, which were 5.3 eV, 5.35 eV, 5.657 eV, and 8.47 eV, depending on the
computational methodology. When Na₂O is introduced into
pure SiO₂, first the band gap decreases sharply, but the
decreasing rate slows down after x = 0.2. The gap of pristine
glass with x = 0.5 was higher than that with x = 0.33, which
may be due to strained regions after structure relaxation in the
x = 0.5 model.

Fifty configurations were chosen from the AIMD production
runs of the CdSe quantum-dot-doped glasses to conduct
further static DFT calculations. The HOMO−LUMO gap
distribution of each composition is shown in Figure 3. The
average values of the gap were 1.95, 2.11, 1.71, and 1.84 eV,
when x changed from 0 (Figure 3a), 0.25 (Figure 3b), 0.33
(Figure 3c), and 0.5 (Figure 3d), respectively. The relationship
between the HOMO−LUMO gap and the ratio of the Na₂O is
nonlinear and different from the relationship between the
HOMO−LUMO gap of the pristine glass and sodium oxygen
contents.

Experimentally fabricated glasses are inhomogeneous in
microscopic regions due to structure fluctuations in such
disordered materials, in addition to the local fluctuations in the
composition and thermal conditions. With our statistical
sampling of configurations from the molecular dynamics
trajectory, we can represent this diversity in structure (and,
as a consequence, physical and chemical properties) despite
having a relatively small simulation box: each configuration in
our calculation can represent a small part of the real glass
sample, representing different possible interfacial environments
of CdSe quantum dots in the glass matrices. The optical
properties of the glass sample are therefore a statistical
representation of these micro regions. Even though the
composition for each configuration of the same system is
identical, we find that the HOMO−LUMO gap still displays

Figure 3. Histograms of HOMO−LUMO gap for CdSe quantum-dot-doped sodium silicate glasses. The composition of the x Na₂O−(1−x)SiO₂
glass matrices: (a) x = 0; (b) x = 0.25; (c) x = 0.33; and (d) x = 0.5. A Gaussian function was used to fit the data, and SD represents the standard
deviation.
important fluctuations, with standard deviations of 0.15, 0.13, 0.29, and 0.12 eV, respectively. These results demonstrate that the interfacial environment of each QD in glass matrices is rather complicated in terms of the electronic structures.

It should be noted that the standard deviation of \( x = 0.33 \) was much higher than other glasses composition, with an asymmetric distribution featuring a long tail at low values. Therefore, it indicates broad distribution in absorption and photoluminescence of CdSe QD formed in the glasses, resulting in low photoluminescence quantum yield of QDs in glasses, which is unfavorable for their potential application. Thus, it gives guidance to the compositional design of glass matrix that \( x = 0.33 \) is not a desirable ratio if we want to fabricate highly luminescent glasses.

To have a better understanding of the impact of glass matrices on the electronic structures of the CdSe QD, the density of states and projected density of states of the pristine glass, pristine QD, and CdSe quantum-dot-doped glass were calculated and are plotted in Figure 4. For the final structure of \( x = 0 \) (Figure 4a), \( x = 0.25 \) (Figure 4b), and \( x = 0.5 \) (Figure 4d), the HOMO orbitals and LUMO orbitals are determined by the hybrid QD. Interestingly, the HOMO orbitals in \( x = 0.33 \) (Figure 4c) appear different compared to other composition, because it is determined by hybrid glass instead of the hybrid QD, while the LUMO orbitals are still decided by hybrid QD. These results demonstrate two different luminescence mechanisms of CdSe quantum-dot-doped glasses, which we depict in Figure 5. For what we will call type I (Figure 5a), when the electrons are excited from the hybrid glass (process i), they will leave holes on the valence band. These photogenerated electrons will relax to the LUMO orbitals formed by the hybrid QD. The holes and electrons are in different spatial locations, resulting in indirect exciton recombination (process ii) and enhanced nonradiative recombination (process iii). Figure 4. Density of states and projected density of states of the final structure of the AIMD calculation of CdSe quantum-dot-doped glasses. The hybrid glass included O atoms, Na atoms, and Si atoms, while hybrid QD was composed of Cd atoms and Se atoms. Cd−O (and Se−Na) represent the projected density of states of the Cd (respectively, Se) atoms bonded with O (Na) atoms and O (Na) atoms bonded with Cd (Se) atoms. The HOMO and the Fermi level are set at 0 eV. (a−d) Compositions of \( x = 0, 0.25, 0.33, \) and 0.5, respectively.

Figure 5. Energy diagrams demonstrating two different luminescence mechanism of CdSe quantum-dot-doped glasses. The HOMO is determined by the hybrid glass in type I (a), while it is decided by the hybrid QD in type II (b).
recombination (process iii). For what we will call type II (Figure 5b), both electrons and holes are generated in hybrid QD upon light excitation (process i), leading to the direct exciton recombination (process ii) and subsequent intrinsic emission from hybrid QD (process iii). On the other hand, the bound electron−hole pairs can also dissociate into free charge carriers. This is a complex process that is influenced by various factors such as shape and size of the QDs, surface defects, and coupling with phonon modes.30 However, in this work, there is no clear boundary between the hybrid QD and hybrid glass. Besides, the difference between the energy of LUMO of pristine glass and that of hybrid QD is big and the exciton binding energy of CdSe QDs (eV) is much bigger than that of bulk semiconductor (meV). Therefore, there is little possibility of forming free electrons and holes in this model.

In addition, we also looked deeper into the impact of interfacial bonds. As we can clearly see from Figure 4, the projected densities of states of Se−Na (magenta solid line) and Cd−O (blue line) contribute to the peak at the bottom of the conduction band. However, the peak formed by the hybrid QD (red solid line), Se−Na, and Cd−O are at the same position. Therefore, we can conclude that both Se−Na and Cd−O linkages contribute to the frontier orbitals but do not shift their energy level.

We further analyzed the density of states of 50 individual configurations of each composition. Based on the DOS, each configuration was assigned to one of the types of luminescence mechanism described above (Figure 6). Importantly, the luminescence mechanisms of CdSe quantum-dot-doped glasses were found to be strongly dependent on the composition of the glass matrices. When $x = 0, 0.25, or 0.5$, the luminescence was mainly determined by the direct transition between HOMO and LUMO of hybrid QD (type I mechanism). On the other hand, when $x = 0.33$, the luminescence was mainly determined by the indirect transition between the LUMO of hybrid QD and HOMO of hybrid glass (type I mechanism), and only 7 out of 50 configurations were classified to type II. Thus, we see again that the composition of the glass with $x = 0.33$ is less favorable for practical applications. For one thing, the standard deviation of $x = 0.33$ (Figure 3) is much higher than other compositions, contributing to the broad distribution of photoluminescence. In addition to that, the electrons and holes are generated in different spatial regions, tending to induce indirect recombination. Both effects are likely to reduce the quantum efficiency.

Traditionally, the luminescence of the CdSe quantum-dot-doped glass can be divided into two competitive parts, which are defect emission and intrinsic emission.5,31 The defect emission is attributed to trap states formed by surface defects such as dangling bond, vacancies, and so on. Meanwhile, the intrinsic emission of CdSe QD is related to the direct recombination of the excitons formed within the QD. However, according to the atomic structure and the HOMO−LUMO gap as well as the density of states of each configuration, we can establish a new luminescence mechanism quite different from the models proposed in previously published work. We hypothesize that the luminescence of CdSe quantum dots embedded in silicate glass matrices observed in experiments was neither defect emission induced by the surface defects nor intrinsic emission from pristine QD as is commonly proposed, but the intrinsic emission of the hybrid QD and hybrid glasses.

It is a highly hybrid system, in which the surface defects and core cannot be separately discussed because their coupling is too important due to strong reorganization at the interface. First, we analyzed the structural properties of the 200 configurations at all compositions, and large-scale structure reconstruction can be observed in each configuration: it is hard for the core to maintain its bulk crystalline structure. Na atoms were found to be bonded with the Se atoms even in the core of
the QD. Therefore, there is little possibility for the CdSe QD to keep a “pristine” state. Second, only a small ratio of configurations was found to have the same value of HOMO−LUMO gap as the pristine QD (Figure 6) with a broad distribution of values instead, further proving the little possibility of the existence of pristine QD in glass matrices. Third, no midgap state was found to be formed inside the HOMO−LUMO gap, as observed in some colloidal quantum dots\(^7\,^9\,^{10}\,^{32}\,^{33}\), where the orbital of pristine QD still exists. Therefore, pristine QD is less likely to exist in the sodium silicate glass. The emission of the CdSe quantum-dot-doped glasses originated from the exciton recombination of the hybrid system formed by the hybrid QD and hybrid glass and cannot be described by simplistic models.

The main goal of this work is to probe the contribution of each component to frontier orbitals and possible structural origins of low quantum efficiency, rather than describing excited states of the modeling system. The inherent limitations of the HOMO−LUMO gap calculation based on the exact Kohn−Sham framework or Kohn−Sham approximations are the discontinuity of the fundamental derivative of exchange-correlation functionals.\(^{34}\,^{35}\,^{36}\,^{37}\) In molecules, the HOMO−LUMO gap obtained by the exact Kohn−Sham potential is still different from the fundamental gap (the sum of optical gap and exciton binding energy), which is theoretically expected to be an approximation to optical gap. However, in solids, the mismatch between the full, local, stabilization by the xc hole for the Kohn−Sham unoccupied state and the weak stabilization by the actual, screened hole that account for the discrepancy between the Kohn−Sham gap and excitation energy.\(^{37}\) Although the HOMO−LUMO gap obtained by Hartree−Fock method is physically close to the fundamental gap, it tends to overestimate the ionization energy and underestimate the electron affinity due to its \(\delta N\)-concave nature of functional, resulting in bigger value of fundamental gap.\(^{38}\)

Hybrid functionals have been proved to predict the fundamental gap accurately for certain systems. The PBE0 functional employed in this work contain both GGA and HF components; thus, the inherent delocalization error of GGA and localization error of HF can cancel but not completely.\(^{36}\) However, in terms of the description of excited states, it has been found that the hybrid functionals not only destroy the attractive feature of Kohn−Sham method that it approximates the lowest excitation energy but also yield the virtual orbitals with the undesirable diffuseness.\(^{37}\) We do not intend to obtain the value comparable to experimental findings but to give a qualitative description of the impact of each component on electronic structures. Moreover, all of the calculations in our work were based on ground states rather than excited states; it is worth conducting further TD-DFT calculations to calculate the electron−hole recombination probability, looking deeper into its optical properties.

We further investigate the impact of the glass matrix (and its composition) on the electronic structure of the CdSe quantum dot (Figure 7) by comparing the density of states of hybrid and pristine QDs. The top of the valence band of the pristine QD is determined by Se atoms, and the bottom of the conduction band is determined by Cd atoms and the Cd atoms not bonded with O atoms. Se represents the Se atoms not bonded with Na atoms. (a−d) Compositions of \(x = 0, 0.25, 0.33,\) and 0.5, respectively, while (e) corresponds to pristine QD.

![Figure 7. Density of states of the final structure of AIMD simulation of CdSe quantum-dot-doped glasses and pristine QD. Cd(O) represents the Cd atoms bonded with O atoms and Se(Na) represents the Se atoms bonded with Na atoms. Cd represents the Cd atoms not bonded with O atoms. Se represents the Se atoms not bonded with Na atoms. (a−d) Compositions of \(x = 0, 0.25, 0.33,\) and 0.5, respectively, while (e) corresponds to pristine QD.](https://doi.org/10.1021/acs.jpcc.1c04665)
Finally, we investigated the effect of the presence of the CdSe QDs on the electronic structure of the glasses. There are two typical segments on the upper valence band of pristine glasses as well as the hybrid glasses (Figure 8). They mainly consist of oxygen 2p lone pair orbitals of the highest orbitals and in bonding states between silicon sp3 hybrids and oxygen 2p orbitals from −10 to −5 eV.41,42 It is apparent that the introduction of CdSe QD sharply decreases the gap from that of the pristine glass, with the bottom of the conduction band shifting to the lower-energy part. Besides, an additional peak appears at the shoulder of the upper segments formed by the oxygen lone pair orbitals in CdSe quantum-dot-doped glasses, which can be attributed to nonbonding oxygen.

Although we identified seven configurations (red points in Figure 6c), out of 50, for which the luminescence mechanism belongs to type II when x = 0.33, the impact of the CdSe QD on the pristine glass and the effect of the pristine glass on CdSe is qualitatively the same as other configurations. For example, from the configuration obtained from the AIMD production run at 7 ps with x = 0.33, it is apparent that the top of the valence band is determined by the hybrid QD, so the luminescence mechanism of this configuration is assigned to type II (Figure 9a), different from the majority of the configurations of x = 0.33. However, the impact of glass or CdSe QD demonstrates similar tendency. A much smaller HOMO–LUMO gap of hybrid glasses compared with pristine glasses (Figure 9b) can be observed, indicating that the presence of CdSe QD can greatly decrease the HOMO–LUMO gap of pristine QD. The hybrid orbitals formed by Se atoms and Na atoms played a dominant role in the HOMO orbitals. Meanwhile, similarly, the HOMO orbital is found to be decided by Se atoms, which are bonded with Na atoms in hybrid QD (Figure 9c).

**CONCLUSIONS AND PERSPECTIVES**

To the best of our knowledge, we present here the first realistic electronic structure calculations of CdSe quantum-dot-doped glasses. By varying the amount of sodium oxygen in x Na2O–(1−x) SiO2 glass matrices, the compositional dependency of structural reconstruction and luminescence mechanisms was illustrated, which is of prominent importance for exploring the optical properties of these materials.

The additional introduction of Na2O in the glass matrix is shown to promote the formation of Cd–O bonds and Se–Na bonds, breaking some Cd–Se bonds. Meanwhile, the luminescence mechanisms show an important dependency on the composition of the glass matrix surrounding the quantum dot. When x = 0, 0.25, and 0.5, the HOMO is decided by the hybrid QD, while the HOMO is determined by the hybrid glasses in most configurations for x = 0.33.

The presence of the CdSe quantum dot in the glass sharply decreases the HOMO–LUMO gap of the interfacial glass...
(near the CdSe QD) compared with pristine glasses. The interfacial bonds contribute greatly to the frontier orbitals, and in particular, we do not observe the formation of any midgap state within the HOMO–LUMO gap. Besides, the HOMO–LUMO gap shows a broader distribution at $x = 0.33$, compared with other compositions. The spatial separation of the holes and electrons in most configurations of $x = 0.33$ and its wide HOMO–LUMO distribution predict a low quantum efficiency of this composition.

In particular, given the atomic structure, HOMO–LUMO gap distribution and density of states of the CdSe quantum-dot-doped glasses, it is apparent that there is little possibility of the existence of the pristine QD, demonstrating no intrinsic emission from the pristine QD. Thus, a totally new energy diagram model is proposed, very different from that known for colloidal QDs, and also from the model generally proposed in the literature for CdSe quantum-dot-doped glasses (based on indirect experimental evidence). The CdSe QD is surrounded by the glass matrix, where the structural reconstruction happens, forming the structure of hybrid QD and hybrid glass. Due to the low concentration of the QDs in glass matrices, these hybrid systems only exist in nanoscale regions, surrounded by dense amorphous pristine glass, acting as a shell toward the hybrid systems composed of hybrid QD and hybrid glasses. Therefore, the emission of the CdSe quantum-dot-doped glasses originates from the intrinsic emission of the hybrid systems, rather than the intrinsic and defect emission from the pristine QD and its trap states in the classical model.

Our results help explain why the quantum efficiency of CdSe quantum-dot-doped glass is much lower than their colloidal counterparts. On the one hand, there is no intrinsic emission from the pristine QD due to the strong impact of the glass matrices on CdSe QD. On the other hand, from the experimental part, the glass is inhomogeneous in micro regions due to the fluctuation of the thermal condition and composition. Even with the same composition of these hybrid systems, the electronic structure of each configuration was variable, exhibiting different optical properties, resulting in the broad emission.

Our results demonstrate the crucial importance of the composition of glass matrix, giving a straightforward illustration of the electronic structure of CdSe quantum-dot-doped glasses, enhancing the mature understanding of its luminescence mechanism from a different perspective. Especially, these luminescence mechanisms can be altered by varying the sodium content in glass matrices. We expect, from the results obtained in our work, the composition of $x = 0.33$ to exhibit poor quantum efficiency. Meanwhile, the composition of $x = 0.5$ is not stable due to the high sodium contents in glass matrices and the very high melting point of $x = 0$. Therefore, the composition of $x = 0.25$ should be an ideal composition for the glass matrices, with holes and electrons confined in the hybrid quantum dot, higher average HOMO–LUMO gap compared with pristine quantum dot, and narrow HOMO–LUMO gap distribution.

### Figure 9
Configuration obtained from the AIMD production run with $x = 0.33$ at 7 ps. 
(a) Density of states of hybrid glass, hybrid QD, pristine QD, pristine glass ($x = 0.33$), Cd–O bonds, and Se–Na bonds. 
(b) Density of states of the hybrid glasses and pristine glasses ($x = 0.33$). 
(c) Density of states of hybrid QD and pristine QD.

### ASSOCIATED CONTENT
Supporting Information
The Supporting Information is available free of charge at https://pubs.acs.org/doi/10.1021/acs.jpcc.1c04665.
Computational details of generation of glass matrix structure and CdSe quantum-dot-doped glasses structure (PDF) 
Geometry structures of the structure used for electronic structures calculations and input files for the classical MD, AIMD, and electronic structure calculation (ZIP)

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Notes

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